



On the Otimes Operator Related to Nonlinear Heat Equation

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Abstract : In this paper, we are finding nonlinear heat equation in n-dimensional. By method of Fourier transform in sense of Distribution theory we obtain the solution in the convolution form. On the suitable we obtain the solution nonlinear triharmonic heat equation.

Keywords : Fourier transform; temper distribution; Diamond operator.

2010 Mathematics Subject Classification : 46F10; 46F12.

1 Introduction

It is well known that for the heat equation

$$\frac{\partial}{\partial t} u(x, t) = c^2 \Delta u(x, t) \quad (1.1)$$

with the initial condition

$$u(x, 0) = f(x)$$

where $\Delta = \sum_{i=1}^n \frac{\partial^2}{\partial x_i^2}$ is the Laplace operator and $(x, t) = (x_1, x_2, \dots, x_n, t) \in \mathbb{R}^n \times (0, \infty)$, we obtain

$$u(x, t) = \frac{1}{(4c^2\pi t)^{n/2}} \int_{\mathbb{R}^n} \exp\left(-\frac{|x-y|^2}{4c^2t}\right) f(y) dy \quad (1.2)$$

as the solution of (1.1). Now, (1.2) can be written $u(x, t) = E(x, t) * f(x)$ where

$$E(x, t) = \frac{1}{(4c^2\pi t)^{n/2}} \exp\left(-\frac{|x|^2}{4c^2t}\right). \quad (1.3)$$

$E(x, t)$ is called *the heat kernel*, where $|x|^2 = x_1^2 + x_2^2 + \dots + x_n^2$ and $t > 0$, [1, pp. 208–209].

Next, Nonlaopon and Kananthai [2] study the equation

$$\frac{\partial}{\partial t} u(x, t) = c^2 \square u(x, t),$$

and \square^k is the ultra-hyperbolic operator iterated k - times and is defined by

$$\square^k = \left(\frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2} + \dots + \frac{\partial^2}{\partial x_p^2} - \frac{\partial^2}{\partial x_{p+1}^2} - \frac{\partial^2}{\partial x_{p+2}^2} - \dots - \frac{\partial^2}{\partial x_{p+q}^2} \right)^k. \quad (1.4)$$

They obtained the ultra-hyperbolic heat kernel

$$E(x, t) = \frac{i^q}{(4c^2\pi t)^{n/2}} \exp\left(-\frac{\sum_{i=1}^p x_i^2 - \sum_{j=p+1}^{p+q} x_j^2}{4c^2t}\right), \quad (1.5)$$

where $p + q = n$ and n is the dimension of the Euclidean space R^n and $i = \sqrt{-1}$ for finding the kernel $E(x, t)$.

In 1996, Kananthai [3] first introduced the operator \diamond^k and is named Diamond operator and is defined by

$$\diamond^k = \left(\left(\sum_{i=1}^p \frac{\partial^2}{\partial x_i^2} \right)^2 - \left(\sum_{j=p+1}^{p+q} \frac{\partial^2}{\partial x_j^2} \right)^2 \right)^k. \quad (1.6)$$

The operator \diamond^k can be written as the product of the operators in the form

$$\diamond^k = \Delta^k \square^k = \square^k \Delta^k, \quad (1.7)$$

where Δ^k is the Laplacian operator iterated k - times and is defined by

$$\Delta^k = \left(\frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2} + \dots + \frac{\partial^2}{\partial x_n^2} \right)^k. \quad (1.8)$$

The Fourier transform of the Diamond operator also has been studied and the elementary solution of such operator, [4]. In 2009, Satsanit [5] has first introduced

the operator \otimes^k , where \otimes^k is defined by

$$\begin{aligned} \otimes^k &= \left(\left(\sum_{i=1}^p \frac{\partial^2}{\partial x_i^2} \right)^3 - \left(\sum_{j=p+1}^{p+q} \frac{\partial^2}{\partial x_j^2} \right)^3 \right)^k \\ &= \left(\sum_{i=1}^p \frac{\partial^2}{\partial x_i^2} - \sum_{j=p+1}^{p+q} \frac{\partial^2}{\partial x_j^2} \right)^k \\ &\quad \times \left[\left(\sum_{i=1}^p \frac{\partial^2}{\partial x_i^2} \right)^2 + \left(\sum_{i=1}^p \frac{\partial^2}{\partial x_i^2} \right) \left(\sum_{j=p+1}^{p+q} \frac{\partial^2}{\partial x_j^2} \right) + \left(\sum_{j=p+1}^{p+q} \frac{\partial^2}{\partial x_j^2} \right)^2 \right]^k \\ &= (\square)^k \left(\Delta^2 - \frac{1}{4}(\Delta + \square)(\Delta - \square) \right)^k \\ &= \left(\frac{3}{4} \diamond \Delta + \frac{1}{4} \square^3 \right)^k, \end{aligned} \tag{1.9}$$

where Δ , \square and \diamond are defined by (1.8), (1.4) and (1.6) with $k = 1$ respectively.

Now, the purpose of this work is to study the equation

$$\frac{\partial}{\partial t} u(x, t) - c^2(-\otimes)^k u(x, t) = f(x, t, u(x, t)) \tag{1.10}$$

for $(x, t) \in \mathbb{R}^n \times (0, \infty)$. We consider the equation (1.10) which is in the form of nonlinear heat equation with the following conditions on u and f as follows

- (1) $u(x, t) \in C^{(6k)}(\mathbb{R}^n)$ for any $t > 0$ where $C^{(6k)}(\mathbb{R}^n)$ is the space of continuous function with $6k$ -derivatives.
- (2) f satisfies the Lipchitz condition,

$$|f(x, t, u) - f(x, t, w)| \leq A|u - w|$$

where A is constant with $0 < A < 1$.

- (3) $\int_0^\infty \int_{\mathbb{R}^n} |t^{-n/6k} f(x, t, u(x, t))| dx dt < \infty$ for $x = (x_1, x_2, \dots, x_n) \in \mathbb{R}^n$, $0 < t < \infty$ and $u(x, t)$ is continuous function on $\mathbb{R}^n \times (0, \infty)$.

Under such conditions of f and u and for the spectrum of $E(x, t)$, we obtain the convolution

$$u(x, t) = E(x, t) * f(x, t, u(x, t))$$

as a unique solution of (1.10) where $E(x, t)$ is defined by (2.9).

2 Preliminaries

Definition 2.1. Let $f(x) \in L_1(\mathbb{R}^n)$ -the space of integrable function in \mathbb{R}^n . The Fourier transform of $f(x)$ is defined by

$$\widehat{f}(\xi) = \frac{1}{(2\pi)^{n/2}} \int_{\mathbb{R}^n} e^{-i(\xi, x)} f(x) dx, \quad (2.1)$$

where $\xi = (\xi_1, \xi_2, \dots, \xi_n)$ and $x = (x_1, x_2, \dots, x_n) \in \mathbb{R}^n$, $(\xi, x) = \xi_1 x_1 + \xi_2 x_2 + \dots + \xi_n x_n$ and $dx = dx_1 dx_2 \dots dx_n$.

Also, the inverse of Fourier transform is defined by

$$f(x) = \frac{1}{(2\pi)^{n/2}} \int_{\mathbb{R}^n} e^{i(\xi, x)} \widehat{f}(\xi) d\xi. \quad (2.2)$$

If f is a distribution with compact supports then Eq.(2.1) can be written as

$$\widehat{f}(\xi) = \frac{1}{(2\pi)^{n/2}} \langle f(x), e^{-i\xi \cdot x} \rangle. \quad (2.3)$$

[6, Theorem 7.4-3, p.187].

Lemma 2.2. Given the function

$$f(x) = \exp \left(\left(\sum_{i=1}^p x_i^2 \right)^3 - \left(\sum_{j=p+1}^{p+q} x_j^2 \right)^3 \right)^k$$

where

$$(x_1, x_2, \dots, x_n) \in \mathbb{R}^n, \quad p + q = n, \quad \left(\sum_{i=1}^p x_i^2 \right)^3 < \left(\sum_{j=p+1}^{p+q} x_j^2 \right)^3$$

and k is the positive odd number. Then

$$\left| \int_{\mathbb{R}^n} f(x) dx \right| \leq \frac{\pi^{n/2}}{9k^2} \cdot \frac{\Gamma(\frac{p}{6k})\Gamma(\frac{q}{6k})}{\Gamma(\frac{p}{2})\Gamma(\frac{q}{2})},$$

where $\frac{p+q}{2} = \frac{n}{2}$ and Γ denoted the gamma function. That is $\int_{\mathbb{R}^n} f(x) dx$ is bounded.

Proof.

$$\begin{aligned} \int_{\mathbb{R}^n} f(x) dx &= \int_{\mathbb{R}^n} \exp \left(\left(\sum_{i=1}^p x_i^2 \right)^3 - \left(\sum_{j=p+1}^{p+q} x_j^2 \right)^3 \right)^k dx \\ &= \int_{\mathbb{R}^n} \exp \left(- \left(\left(\sum_{j=p+1}^{p+q} x_j^2 \right)^3 - \left(\sum_{i=1}^p x_i^2 \right)^3 \right)^k \right) dx \end{aligned}$$

for k is a positive odd number. By changing the coordinate, we put

$$\begin{aligned} x_1 &= iy_1, \quad x_2 = iy_2, \dots, \quad x_p = iy_p \\ dx_1 &= idy_1, \quad dx_2 = idy_2, \dots, \quad dx_p = idy_p \end{aligned}$$

and

$$\begin{aligned} x_{p+1} &= y_{p+1}, \quad x_{p+2} = y_{p+2}, \dots, \quad x_{p+q} = y_{p+q} \\ dx_{p+1} &= dy_{p+1}, \quad dx_{p+2} = dy_{p+2}, \dots, \quad dx_{p+q} = dy_{p+q}, \end{aligned}$$

where $i = \sqrt{-1}$. Then we obtain

$$\int_{\mathbb{R}^n} f(x)dx = i^p \int_{\mathbb{R}^n} \exp \left(- \left(\left(\sum_{j=p+1}^{p+q} y_j^2 \right)^3 + \left(\sum_{i=1}^p y_i^2 \right)^3 \right) \right)^k dy \quad (2.4)$$

Let us transform to bipolar coordinates defined by

$$y_1 = rw_1, y_2 = rw_2, \dots, \quad y_p = rw_p$$

$$\text{and } y_{p+1} = sw_{p+1}, y_{p+2} = sw_{p+2}, \dots, y_{p+q} = sw_{p+q}, \quad p + q = n$$

where $w_1^2 + w_2^2 + \dots + w_p^2 = 1$ and $w_{p+1}^2 + w_{p+2}^2 + \dots + w_{p+q}^2 = 1$, Thus

$$\int_{\mathbb{R}^n} f(x)dx = i^p \int_{\mathbb{R}^n} \exp \left(- (r^6 + s^6)^k \right) r^{p-1} s^{q-1} dr ds d\Omega_p d\Omega_q \quad (2.5)$$

where

$$dy = r^{p-1} s^{q-1} dr ds d\Omega_p d\Omega_q, \quad (2.6)$$

$d\Omega_p$ and $d\Omega_q$ are the elements of surface area on the unit sphere in R^p and R^q respectively. By computing directly, we obtain

$$\int_{\mathbb{R}^n} f(x)dx = i^p \Omega_p \Omega_q \int_0^\infty \int_0^\infty \exp \left(- (r^6 + s^6)^k \right) r^{p-1} s^{q-1} dr ds \quad (2.7)$$

where $\Omega_p = \frac{2\pi^{p/2}}{\Gamma(p/2)}$ and $\Omega_q = \frac{2\pi^{q/2}}{\Gamma(q/2)}$. Since $(r^6 + s^6)^k \geq r^{6k} + s^{6k}$, we have

$$\exp \left(- (r^6 + s^6)^k \right) \leq \exp \left(- (r^{6k} + s^{6k}) \right)$$

Thus

$$\begin{aligned} \left| \int_{\mathbb{R}^n} f(x)dx \right| &\leq \Omega_p \Omega_q \int_0^\infty \int_0^\infty \exp \left(- r^{6k} - s^{6k} \right) r^{p-1} s^{q-1} dr ds \\ &= \Omega_p \Omega_q \int_0^\infty \exp \left(- r^{6k} \right) r^{p-1} dr \int_0^\infty \exp \left(- s^{6k} \right) s^{q-1} ds. \end{aligned}$$

Put $u = r^{6k}$, $dr = \frac{1}{6k}u^{\frac{1}{6k}-1}du$ and $v = s^{6k}$, $ds = \frac{1}{6k}v^{\frac{1}{6k}-1}dv$ in the above equation, we obtain

$$\begin{aligned} \left| \int_{\mathbb{R}^n} f(x)dx \right| &\leq \frac{\Omega_p \Omega_q}{(6k)^2} \int_0^\infty e^{-u} u^{\frac{1}{6k}-1} du \int_0^\infty e^{-v} v^{\frac{1}{6k}-1} dv \\ &= \frac{\Omega_p \Omega_q}{(6k)^2} \Gamma\left(\frac{p}{6k}\right) \Gamma\left(\frac{q}{6k}\right) \\ &= \frac{2\pi^{p/2} 2\pi^{q/2}}{(6k)^2} \frac{\Gamma(\frac{p}{6k}) \Gamma(\frac{q}{6k})}{\Gamma(\frac{p}{2}) \Gamma(\frac{q}{2})} \\ &= \frac{\pi^{n/2}}{9k^2} \frac{\Gamma(\frac{p}{6k}) \Gamma(\frac{q}{6k})}{\Gamma(\frac{p}{2}) \Gamma(\frac{q}{2})}, \end{aligned}$$

where $\frac{p+q}{2} = \frac{n}{2}$. That is $\int_{\mathbb{R}^n} f(x)dx$ is bounded. □

Lemma 2.3. (The Fourier transform of $\otimes^k \delta$)

$$\mathcal{F} \otimes^k \delta = \frac{(-1)^{3k}}{(2\pi)^{n/2}} \left[(\xi_1^2 + \xi_2^2 + \dots + \xi_p^2)^3 - (\xi_{p+1}^2 + \xi_{p+2}^2 + \dots + \xi_{p+q}^2)^3 \right]^k,$$

where \mathcal{F} is the Fourier transform defined by (2.1) and if the norm of ξ is given by $\|\xi\| = (\xi_1^2 + \xi_2^2 + \dots + \xi_n^2)^{1/2}$ then

$$\mathcal{F} \otimes^k \delta \leq \frac{M}{(2\pi)^{n/2}} \|\xi\|^{6k}.$$

Since M is positive constant thus $\mathcal{F} \otimes^k \delta$ is bounded and continuous on the space \mathcal{S}' of the tempered distribution. Moreover, by Eq.(2.9)

$$\otimes^k \delta = \mathcal{F}^{-1} \frac{(-1)^{3k}}{(2\pi)^{n/2}} \left[(\xi_1^2 + \xi_2^2 + \dots + \xi_p^2)^3 - (\xi_{p+1}^2 + \xi_{p+2}^2 + \dots + \xi_{p+q}^2)^3 \right]^k$$

Proof. See [5]. □

Lemma 2.4. Given the operator

$$L = \frac{\partial}{\partial t} - c^2(-\otimes)^k, \tag{2.8}$$

where $(-\otimes)^k$ is the operator defined by (1.5), k is a positive odd integer, $u(x, t)$ is an unknown function for $(x, t) = (x_1, x_2, \dots, x_n, t) \in \mathbb{R}^n \times (0, \infty)$, and c is a positive constant. Then we obtain

$$E(x, t) = \frac{1}{(2\pi)^n} \int_{\mathbb{R}^n} \exp \left[-c^2 \left[\left(\sum_{j=p+1}^{p+q} \xi_j^2 \right)^3 - \left(\sum_{i=1}^p \xi_i^2 \right)^3 \right] t + i(\xi, x) \right] d\xi, \tag{2.9}$$

where $\sum_{j=p+1}^{p+q} \xi_j^2 > \sum_{i=1}^p \xi_i^2$, as an elementary solution for the operator L defined by (2.8).

Proof. We have to find function $E(x, t)$ from the equation

$$L(E(x, t)) = \delta(x, t),$$

where $\delta(x, t)$ is dirac delta function for $(x, t) \in \mathbb{R}^n \times (0, \infty)$. We can write

$$\frac{\partial}{\partial t} E(x, t) - c^2(-\otimes)^k E(x, t) = \delta(x) \cdot \delta(t). \tag{2.10}$$

By taking the Fourier transform defined by (2.1) to both sides of (2.10), we obtain

$$\frac{\partial}{\partial t} \widehat{E}(\xi, t) - c^2 \left[\left(\sum_{i=1}^p \xi_i^2 \right)^3 - \left(\sum_{j=p+1}^{p+q} \xi_j^2 \right)^3 \right]^k \widehat{E}(\xi, t) = \frac{1}{(2\pi)^{n/2}} \delta(t),$$

which has solution

$$\widehat{E}(\xi, t) = \frac{H(t)}{(2\pi)^{n/2}} \exp \left[c^2 t \left(\left(\sum_{i=1}^p \xi_i^2 \right)^3 - \left(\sum_{j=p+1}^{p+q} \xi_j^2 \right)^3 \right)^k \right] \tag{2.11}$$

where $H(t)$ is a heaviside function and $H(t) = 1$ for $t > 0$. Since k is a positive odd number and $\sum_{j=p+1}^{p+q} \xi_j^2 > \sum_{i=1}^p \xi_i^2$, thus $\widehat{E}(\xi, t)$ is bounded and can be written by

$$\widehat{E}(\xi, t) = \frac{H(t)}{(2\pi)^{n/2}} \exp \left[-c^2 t \left(\left(\sum_{j=p+1}^{p+q} \xi_j^2 \right)^3 - \left(\sum_{i=1}^p \xi_i^2 \right)^3 \right)^k \right]. \tag{2.12}$$

Now, the inverse Fourier transform

$$E(x, t) = \frac{1}{(2\pi)^{n/2}} \int_{\mathbb{R}^n} e^{i(\xi, x)} \widehat{E}(\xi, t) d\xi.$$

Thus

$$\begin{aligned} E(x, t) &= \frac{H(t)}{(2\pi)^{n/2}} \int_{\mathbb{R}^n} e^{i(\xi, x)} \exp \left[-c^2 t \left(\left(\sum_{j=p+1}^{p+q} \xi_j^2 \right)^3 - \left(\sum_{i=1}^p \xi_i^2 \right)^3 \right)^k \right] \\ &= \frac{1}{(2\pi)^{n/2}} \int_{\mathbb{R}^n} e^{i(\xi, x)} \exp \left[-c^2 t \left(\left(\sum_{j=p+1}^{p+q} \xi_j^2 \right)^3 - \left(\sum_{i=1}^p \xi_i^2 \right)^3 \right)^k \right] \end{aligned}$$

or

$$E(x, t) = \frac{1}{(2\pi)^n} \int_{\mathbb{R}^n} \exp \left[-c^2 t \left[\left(\sum_{j=p+1}^{p+q} \xi_j^2 \right)^3 - \left(\sum_{i=1}^p \xi_i^2 \right)^3 \right]^k + i(\xi, x) \right] d\xi. \tag{2.13}$$

as an elementary solution of (2.10) and bounded if $\sum_{j=p+1}^{p+q} \xi_j^2 > \sum_{i=1}^p \xi_i^2$ and k is positive odd number. \square

3 Main Results

Theorem 3.1. *The kernel $E(x, t)$ defined by (2.13) has the following properties:*

- (1) $E(x, t) \in C^\infty$ -the space of continuous function for $x \in \mathbb{R}^n, t > 0$ with infinitely differentiable.
- (2) $\left(\frac{\partial}{\partial t} - c^2(-\otimes)^k \right) E(x, t) = 0$ for $t > 0$ and k is positive odd number.
- (3) $E(x, t) > 0$ for $t > 0$.
- (4) $|E(x, t)| \leq \frac{M(t)}{9.2^n \pi^{n/2} k^2 (c^2 t)^{\frac{n}{6k}}} \frac{\Gamma(\frac{p}{6k}) \Gamma(\frac{q}{6k})}{\Gamma(\frac{p}{2}) \Gamma(\frac{q}{2})}$, for $t > 0$, where $M(t)$ is a positive constant. Thus $E(x, t)$ is bounded for any fixed $t > 0$.
- (5) $\lim_{t \rightarrow 0} E(x, t) = \delta$.

Proof. (1) From (2.13), since

$$\frac{\partial^n}{\partial x^n} E(x, t) = \frac{1}{(2\pi)^n} \int_{\mathbb{R}^n} \frac{\partial^n}{\partial x^n} \exp \left[c^2 \left[\left(\sum_{j=p+1}^{p+q} \xi_j^2 \right)^3 - \left(\sum_{i=1}^p \xi_i^2 \right)^3 \right]^k t + i(\xi, x) \right] d\xi.$$

Thus $E(x, t) \in C^\infty$ for $x \in \mathbb{R}^n, t > 0$.

(2) By computing directly, we obtain

$$\left(\frac{\partial}{\partial t} - c^2(-\otimes)^k \right) E(x, t) = 0.$$

for $t > 0$ where $E(x, t)$ is defined by (2.13).

(3) $E(x, t) > 0$ for $t > 0$ is obvious by (2.13).

(4) We have

$$E(x, t) = \frac{1}{(2\pi)^n} \int_{\mathbb{R}^n} \exp \left[-c^2 t \left[\left(\sum_{j=p+1}^{p+q} \xi_j^2 \right)^3 - \left(\sum_{i=1}^p \xi_i^2 \right)^3 \right]^k + i(\xi, x) \right] d\xi.$$

put

$$\begin{aligned} \xi_1 &= iy_1, \quad \xi_2 = iy_2, \dots, \quad \xi_p = iy_p \\ d\xi_1 &= idy_1, \quad d\xi_2 = idy_2, \dots, \quad d\xi_p = idy_p \end{aligned}$$

and

$$\begin{aligned} \xi_{p+1} &= y_{p+1}, \quad \xi_{p+2} = y_{p+2}, \dots, \quad \xi_{p+q} = y_{p+q} \\ d\xi_{p+1} &= dy_{p+1}, \quad d\xi_{p+2} = dy_{p+2}, \dots, \quad d\xi_{p+q} = dy_{p+q}, \end{aligned}$$

where $i = \sqrt{-1}$. Thus, we obtain

$$\begin{aligned} E(x, t) &= \frac{i^p}{(2\pi)^n} \int_{\mathbb{R}^n} \exp \left[-c^2 t \left[\left(\sum_{j=p+1}^{p+q} \xi_j^2 \right)^3 - \left(\sum_{i=1}^p \xi_i^2 \right)^3 \right]^k - \sum_{r=1}^p x_r y_r + i \sum_{j=p+1}^{p+q} x_j y_j \right] dy. \\ |E(x, t)| &\leq \frac{M}{(2\pi)^n} \int_{\mathbb{R}^n} \exp \left[-c^2 t \left(\left(\sum_{j=p+1}^{p+q} y_j^2 \right)^3 + \left(\sum_{r=1}^p y_r^2 \right)^3 \right) \right] dy, \quad (3.1) \end{aligned}$$

where M is a positive constant. The same process as Lemma 2.1 then (3.1) becomes

$$|E(x, t)| \leq \frac{M(t)}{9 \cdot 2^n \pi^{n/2} k^2 (c^2 t)^{\frac{n}{6k}}} \frac{\Gamma(\frac{p}{6k}) \Gamma(\frac{q}{6k})}{\Gamma(\frac{p}{2}) \Gamma(\frac{q}{2})}$$

(5) We have

$$E(x, t) = \frac{1}{(2\pi)^n} \int_{\mathbb{R}^n} \exp \left[-c^2 t \left[\left(\sum_{j=p+1}^{p+q} \xi_j^2 \right)^3 - \left(\sum_{i=1}^p \xi_i^2 \right)^3 \right]^k + i(\xi, x) \right] d\xi.$$

Since $E(x, t)$ exists, we have

$$\begin{aligned} \lim_{t \rightarrow 0} E(x, t) &= \frac{1}{(2\pi)^n} \int_{\mathbb{R}^n} e^{i(\xi, x)} d\xi \\ &= \frac{1}{(2\pi)^n} \int_{\mathbb{R}^n} e^{i(\xi, x)} d\xi \\ &= \delta(x), \quad \text{for } x \in \mathbb{R}^n \end{aligned} \quad (3.2)$$

[7, p.396, Eq.(10.2.19b)]. □

Theorem 3.2. *Given the nonlinear equation*

$$\frac{\partial}{\partial t} u(x, t) - c^2(-\otimes)^k u(x, t) = f(x, t, u(x, t)) \quad (3.3)$$

for $(x, t) \in \mathbb{R}^n \times (0, \infty)$. We consider the equation (3.3) which is in the form of nonlinear heat equation with the following conditions on u and f as follows

- (1) $u(x, t) \in C^{(6k)}(\mathbb{R}^n)$ for any $t > 0$ where $C^{(6k)}(\mathbb{R}^n)$ is the space of continuous function with $6k$ -derivatives.
- (2) f satisfies the Lipchitz condition,

$$|f(x, t, u) - f(x, t, w)| \leq A|u - w|$$

where A is constant with $0 < A < 1$.

- (3) $\int_0^\infty \int_{\mathbb{R}^n} |t^{-n/6k} f(x, t, u(x, t))| dxdt < \infty$ for $x = (x_1, x_2, \dots, x_n) \in \mathbb{R}^n$, $0 < t < \infty$ and $u(x, t)$ is continuous function on $\mathbb{R}^n \times (0, \infty)$.

Under such conditions of f and u , we obtain the convolution

$$u(x, t) = E(x, t) * f(x, t, u(x, t))$$

as a unique solution of (3.3) where $E(x, t)$ is an elementary solution defined by (2.9). In particular, if we put $k = 1$ and $p = 0$ in (3.3) then (3.3) reduces to the equation

$$\frac{\partial}{\partial t} u(x, t) - c^2 \Delta^3 u(x, t) = f(x, t, u(x, t)),$$

where is related to the nonlinear triharmonic heat equation.

Proof. Convoluting both sides of (3.1) with $E(x, t)$, that is

$$E(x, t) * \left[\frac{\partial}{\partial t} u(x, t) - c^2 (-\otimes)^k u(x, t) \right] = E(x, t) * f(x, t, u(x, t))$$

or

$$\left[\frac{\partial}{\partial t} E(x, t) - c^2 (-\otimes)^k E(x, t) \right] * u(x, t) = E(x, t) * f(x, t, u(x, t)),$$

so

$$\delta(x, t) * u(x, t) = E(x, t) * f(x, t, u(x, t)).$$

Thus

$$\begin{aligned} u(x, t) &= E(x, t) * f(x, t, u(x, t)) \\ &= \int_{-\infty}^\infty \int_{\mathbb{R}^n} E(r, s) f(x - r, t - s, u(x - r, t - s)) drds \end{aligned}$$

where $E(r, s)$ is given by definition (2.4). We next show that $u(x, t)$ is bounded on $\mathbb{R}^n \times (0, \infty)$. We have

$$\begin{aligned} |u(x, t)| &\leq \int_{-\infty}^\infty \int_{\mathbb{R}^n} |E(r, s)| \cdot |f(x - r, t - s, u(x - r, t - s))| drds \\ &\leq \frac{M}{9 \cdot 2^n \pi^{n/2} k^2 c^{n/3k}} \int_{-\infty}^\infty \int_{\mathbb{R}^n} s^{-n/6k} |f(x - r, t - s, u(x - r, t - s))| drds \\ &\leq \frac{MN}{9 \cdot 2^n \pi^{n/2} k^2 c^{n/3k}} \end{aligned} \tag{3.4}$$

where $N = \int_{-\infty}^{\infty} \int_{\mathbb{R}^n} |f(x-r, t-s, u(x-r, t-s))| dr ds$. Thus $u(x, t)$ is bounded on $\mathbb{R}^n \times (0, \infty)$. To show that $u(x, t)$ is unique. Now, we next to show that $u(x, t)$ is unique. Let $w(x, t)$ be another solution of (3.1), then

$$w(x, t) = E(x, t) * f(x, t, w(x, t))$$

for $(x, t) \in \Omega_0 \times (0, T]$ the compact subset of $\mathbb{R}^n \times [0, \infty)$ and $E(x, t)$ is defined by (2.6).

Now, define $\|u(x, t)\| = \sup_{\substack{x \in \Omega_0 \\ 0 < t \leq T}} |u(x, t)|$.

Now,

$$\begin{aligned} |u(x, t) - w(x, t)| &= |E(x, t) * f(x, t, u(x, t)) - E(x, t) * f(x, t, w(x, t))| \\ &\leq \int_{-\infty}^{\infty} \int_{\mathbb{R}^n} |E(r, s)| \cdot |f(x-r, t-s, u(x-r, t-s)) \\ &\quad - f(x-r, t-s, w(x-r, t-s))| dr ds \\ &\leq A |E(r, s)| \int_{-\infty}^{\infty} \int_{\mathbb{R}^n} |u(x-r, t-s) - w(x-r, t-s)| dr ds \end{aligned}$$

by (2.9) and the condition (2) of the theorem. Now, for $(x, t) \in \Omega_0 \times (0, T]$ we have

$$\begin{aligned} |u - w| &\leq A |E(r, s)| \|u - w\| \int_0^T ds \int_{\Omega_0} dr \\ &= A |E(r, s)| TV(\Omega_0) \|u - w\| \end{aligned} \tag{3.5}$$

where $V(\Omega_0)$ is the volume of the surface on Ω_0 .

Choose $A |E(r, s)| TV(\Omega_0) \leq 1$ or $A \leq \frac{1}{|E(r, s)| TV(\Omega_0)}$. Thus from (3.5),

$$\|u - w\| \leq \alpha \|u - w\| \quad \text{where } \alpha = A |E(r, s)| TV(\Omega_0) \leq 1.$$

It follows that $\|u - w\| = 0$, thus $u = w$. That is the solution u of (3.3) is unique. In particular, if we put $k = 1$ and $p = 0$ in (3.3), then (3.3) reduces to the nonlinear heat equation

$$\frac{\partial}{\partial t} u(x, t) - c^2 \Delta^3 u(x, t) = f(x, t, u(x, t))$$

which has solution

$$u(x, t) = E(x, t) * f(x, t, u(x, t))$$

where $E(x, t)$ is defined by (2.9) with $k = 1$ and $p = 0$. That is complete of proof. □

Acknowledgement(s) : The author would like to thank The Thailand Research Fund, The Commission an Higher Education and Graduate School, Maejo University, Chiang Mai, Thailand for financial support and also Prof. Amnuay Kananthai Department of Mathematics, Chiang Mai University for the helpful of discussion.

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(Received 4 August 2012)

(Accepted 4 June 2013)