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# Option Pricing under a Mean Reverting Process with Jump-Diffusion and Jump Stochastic Volatility

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**Abstract**: An alternative option pricing model is proposed, in which the asset prices follow the jump-diffusion and exhibits mean reversion. The stochastic volatility follows the jump-diffusion with mean reversion. We find a formulation for the European-style option in terms of characteristic functions.

**Keywords :** jump-diffusion model; stochastic volatility; characteristic function; option pricing; mean reverting.

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## 1 Introduction

Let  $(\Omega, \mathcal{F}, \mathbb{P})$  be a probability space with filtration  $\mathcal{F} = (\mathcal{F}_t)_{0 \leq t \leq T}$ . All processes that we shall consider in this section will be defined in this space. An asset price model with stochastic volatility has been defined by Heston [1] which has the following dynamics:

$$dS_t = S_t(\mu dt + \sqrt{v_t} dW_t^S), \tag{1.1}$$

$$dv_t = \kappa(\theta - v_t)dt + \sigma\sqrt{v_t}dW_t^v, \qquad (1.2)$$

where  $S_t$  is the asset price,  $\mu \in \Re$  is the rate of return of the asset price,  $v_t$  is the volatility of asset returns,  $\kappa > 0$  is the rate at which the volatility reverts

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toward its long-term mean,  $\theta \in \Re$  is the mean long-term volatility,  $\sigma > 0$  is the volatility of the volatility process,  $W_t^S$  and  $W_t^v$  are standard Brownian motions corresponding to the processes  $S_t$  and  $v_t$ , respectively, with constant correlation  $\rho$ . Bate [2] introduced the jump-diffusion stochastic volatility model by adding log normal jump  $Y_t$  to the Heston stochastic volatility model. In the original formulation of Bate, the model has the following form:

$$dS_{t} = S_{t}(\mu dt + \sqrt{v_{t}} dW_{t}^{S}) + S_{t-}Y_{t} dN_{t}^{S}, \qquad (1.3)$$

where  $N_t^S$  is the Poisson process which corresponds to the underlying asset  $S_t$ ,  $Y_t$  is a proportion of jump size of the asset price (1.1) with log normal distribution and  $S_t$  – means that there is a jump in the value of the process before the jump is used on the left-hand side of the formula. Eraker et al. [3] extended Bate's work by incorporating jumps into the volatility model, i.e.

$$dv_t = \kappa(\theta - v_t)dt + \sigma\sqrt{v_t}dW_t^v + Z_t dN_t^v \tag{1.4}$$

Eraker et al. [3] developed a likelihood-based estimation strategy and provided estimates of parameters, spot volatility, jump times, and jump sizes using S&P 500 and Nasdaq 100 index returns. Moreover, they examined the volatility structure of the S&P and Nasdaq indices and indicated that models with jumps in volatility are preferred over those without jumps in volatility. But they did not provide a closed-form formula for the price of a European call option.

Empirical evidence on mean reversion in financial assets has been produced by Cecchetti et al. [4] and Bessembinder et al. [5], respectively. It has been documented that currency exchange rates also exhibit mean reversion. Jorion and Sweeney [6] show how the real exchange rates revert to their mean levels and Sweeney [7] provides empirical evidence of mean reversion in G-10 nominal exchange rates. Mean reversion also appears in some stock prices as evidenced by Poterba and Summers [8].

In this paper, we consider the problem of finding a closed-form formula for a European call option where the asset price follows mean reverting jump-diffusion and the stochastic volatility with jump.

The rest of this paper is organized as follows. In Section 2, we briefly discuss model descriptions for option pricing. Deriving a formula for a characteristic function is presented in Section 3. Finally, a closed-form formula for a European call option in terms of characteristic functions is presented.

#### 2 Model Descriptions

It is assumed that a risk-neutral probability measure  $\mathcal{M}$  exists. The asset price  $S_t$  under this measure follows a mean reverting jump-diffusion process, and the volatility  $v_t$  follows mean reverting with jump, i.e. our models are governed

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by the following dynamics:

$$dS_t = b\left(a - \ln S_t - \frac{\lambda^S m}{b}\right)S_t dt + \sqrt{v_t}S_t dW_t^S + S_{t-}Y_t dN_t^S$$
(2.1)

$$dv_t = \kappa(\theta - v_t)dt + \sigma\sqrt{v_t}dW_t^v + Z_t dN_t^v$$
(2.2)

where  $S_t, v_t, \kappa, \theta, \sigma, W_t^S$  and  $W_t^v$  are defined as above,  $a \in \Re$  is the mean of longterm asset price return, b > 0 is the rate at which the asset price return reverts toward its long-term mean,  $N_t^S$  and  $N_t^v$  are independent Poisson processes with constant intensities  $\lambda^S$  and  $\lambda^v$  respectively.  $Y_t$  and  $Z_t$  are proportional jump sizes of the asset price (2.1) and the jump size of the volatility process (2.2) respectively. Suppose that  $Y_t$  and  $Z_t$  are independent and identically distributed sequences with densities  $\phi_{Y_t}(y) := \phi_Y(y), \phi_{Z_t}(z) := \phi_Z(z)$  and  $EY_t := m$ . Moreover, we assume that the jump processes  $N_t^S$  and  $N_t^v$  are independent of standard Brownian motions  $W_t^S$  and  $W_t^v$ .

Assume that the asset price  $S_t$  and the volatility  $v_t$  satisfy equations (2.1) and (2.2) respectively. Let  $L_t = \ln S_t$ , by the jump-diffusion chain rule,  $\ln S_t$  satisfies the SDE

$$dL_t = b\left(a - L_t - \frac{\lambda^S m}{b} - \frac{v_t}{2b}\right)dt + \sqrt{v_t}dW_t^S + \ln(1 + Y_t)dN_t^S.$$
 (2.3)

#### **3** Characteristic Functions

We denote the characteristic function for  $L_T = \ln S_T$  as

$$f(x:t,l,v) = E_{\mathcal{M}}[e^{ixL_T}|L_t = l, v_t = v]$$
(3.1)

where  $0 \le t \le T$  and  $i = \sqrt{-1}$ . Here  $L_t$  is the mean reverting asset price process with jumps specified by (2.3) and  $v_t$  is the volatility process specified by (2.2). The generalized Feynman-Kac theorem [9] implies that f(x : t, l, v) solves the following partial integro-differential equation (PIDE):

$$0 = \frac{\partial f}{\partial t} + b\left(a - l - \frac{\lambda^S m}{b} - \frac{v}{2b}\right)\frac{\partial f}{\partial l} + \kappa(\theta - v)\frac{\partial f}{\partial v} + \frac{1}{2}v\frac{\partial^2 f}{\partial l^2} + \rho\sigma v\frac{\partial^2 f}{\partial l\partial v} + \frac{1}{2}\sigma^2 v\frac{\partial^2 f}{\partial v^2} + \lambda^S \int_{\Re} [f(x:t,l+y,v) - f(x:t,l,v)]\phi_Y(y)dy + \lambda^v \int_{\Re} [f(x:t,l,v+z) - f(x:t,l,v)]\phi_Z(z)dz.$$
(3.2)

**Lemma 3.1.** Suppose that  $L_t$  follows the dynamics in (2.3). Then the characteristic function for  $L_T$  can be written in the form

$$f(x:t,l,v) = \exp[B(t,T) + C(t,T)l + D(t,T)v + ixl],$$
(3.3)

where

$$\begin{split} B(t,T) &= (\frac{\lambda^{S}m}{b} - a)ix(e^{-b(T-t)} - 1) - \theta\kappa \int_{t}^{T} D(s,T)ds \\ &+ (T-t)\lambda^{S} \int_{\Re} \left[e^{ixy} - 1\right]\phi_{Y}(y)dy \\ &+ (T-t)\lambda^{v} \int_{\Re} \left[e^{zD(t,T)} - 1\right]\phi_{Z}(z)dz, \\ C(t,T) &= ix(e^{-b(T-t)} - 1), \\ D(t,T) &= U(e^{-b(T-t)}) + \frac{e^{-\kappa(T-t)}V(e^{-b(T-t)})}{-\frac{1}{U(1)} + \frac{\sigma^{2}}{2b} \int_{1}^{e^{-b(T-t)}} h^{\frac{\kappa}{b}} - 1V(h)dh, \end{split}$$

$$\begin{split} U(h) &= \frac{2bh}{\sigma^2} \frac{(\sqrt{1-\rho^2}-\rho i)\frac{\sigma x}{2b} \Phi(a^*,b^*,\frac{h}{\zeta}) + \frac{a}{b^*\zeta} \Phi(a^*+1,b^*+1,\frac{h}{\zeta})}{\Phi(a^*,b^*,\frac{h}{\zeta})},\\ V(h) &= \frac{\Phi^2(a^*,b^*,\frac{1}{\zeta})e^{(\sqrt{1-\rho^2})\frac{\sigma x}{b}(1-h)}}{\Phi^2(a^*,b^*,\frac{h}{\zeta})},\\ h &= e^{-b(T-t)},\\ a^* &= \frac{\frac{b^*}{2}(\sqrt{\rho^2-1}+\rho) + \frac{\sigma}{4b}}{\sqrt{\rho^2-1}},\\ b^* &= 1 - \frac{\kappa}{b},\\ \zeta &= \frac{-b}{\sigma x\sqrt{1-\rho^2}}, \end{split}$$

and  $\Phi(\cdot,\cdot,\cdot)$  is the degenerated hypergeometric function.

*Proof.* From (3.1), it is clear that

$$f(x:T,l,v) = e^{ixl} \tag{3.4}$$

which is the boundary condition of PIDE (3.2). This implies that

$$B(T,T) = C(T,T) = D(T,T) = 0.$$
(3.5)

Substituting (3.3) in (3.2) and using the fact that the function f is never zero, we obtain

$$0 = [B_t + (ba - \lambda^S m)(C + ix) + \kappa \theta D + \lambda^S \int_{\Re} [e^{ixy} - 1] \phi_Y(y) dy + \lambda^v \int_{\Re} [e^{zD} - 1] \phi_Z(z) dz] + [C_t - b(C + ix)] l + [D_t + \frac{1}{2}(C + ix) + \frac{1}{2}(C + ix)^2 - \kappa D + \frac{1}{2}\sigma^2 D^2 + \rho\sigma(C + ix)D] v$$
(3.6)

where  $B_t, C_t$  and  $D_t$  are the partial derivatives with respect to t of functions B, C and D respectively.

This reduces the problem to one of solving three, much simpler, ordinary differential equations:

$$B_t + (ba - \lambda^S m)(C + ix) + \kappa \theta D + \lambda^S \int_{\Re} [e^{ixy} - 1]\phi_Y(y)dy$$
$$+ \lambda^v \int_{\Re} [e^{zD} - 1]\phi_Z(z)dz = 0$$
(3.7)

$$C_t - b(C + ix) = 0$$
 (3.8)

$$D_t + \frac{1}{2}(C+ix)(C+ix-1) - \kappa D + \frac{1}{2}\sigma^2 D^2 + \rho\sigma(C+ix)D = 0$$
(3.9)

subject to boundary conditions (3.5).

The solution to equation (3.8) with the boundary condition C(T,T) = 0 is given by

$$C(t,T) = ix(e^{-b(T-t)} - 1).$$
(3.10)

We now consider equation (3.9). Substituting (3.10) in (3.9), one gets

$$D_t + \frac{1}{2} \left[ ixe^{-b(T-t)} \right] \left[ ixe^{-b(T-t)} - 1 \right] - \kappa D + \frac{1}{2}\sigma^2 D^2 + \rho\sigma ixDe^{-b(T-t)} = 0.$$

Hence,

$$D_t = -\frac{1}{2}\sigma^2 D^2 + \left[\kappa - \rho\sigma ixe^{-b(T-t)}\right] D + \frac{1}{2} \left[x^2 e^{-2b(T-t)} + ixe^{-b(T-t)}\right].$$
 (3.11)

Let  $h = e^{-b(T-t)}$  and we define a new function  $\hat{D}(h(t), T) := D(t, T)$ . Then

$$\frac{\partial D(t,T)}{\partial t} = \frac{\partial D(h,T)}{\partial h} \frac{\partial h}{\partial t}$$
$$= be^{-b(T-t)} \frac{\partial \hat{D}(h,T)}{\partial h}.$$
(3.12)

Substituting (3.12) into (3.11), we obtain the following Riccati equation

$$\frac{\partial \hat{D}}{\partial h} = -\frac{1}{2bh}\sigma^2 \hat{D}^2 + \left(\frac{\kappa}{bh} - \frac{\rho\sigma ix}{b}\right)\hat{D} + \frac{1}{2b}\left(x^2h + ix\right).$$
(3.13)

We shall solve the second order ODE (3.13) together with the initial condition  $\hat{D}(1,T) = 0$ . Let

$$\hat{D}(h,T) = \frac{2bhw'(h)}{\sigma^2 w(h)}$$
(3.14)

and taking the derivative of (3.14) with respect to h, one gets

$$\frac{\partial \hat{D}}{\partial h} = \left[\sigma^2 w(h) \frac{\partial}{\partial h} (2bhw'(h)) - 2bhw'(h) \frac{\partial}{\partial h} (\sigma^2 w(h))\right] \frac{1}{\sigma^4 w^2(h)} \\
= \left[\sigma^2 w(h) \left[2bw'(h) + 2bhw''(h)\right] - 2bh\sigma^2 (w'(h))^2\right] \frac{1}{\sigma^4 w^2(h)}.$$
(3.15)

Substituting (3.14) and (3.15) into (3.13), we have

$$hw''(h) - \left[\left(\frac{\kappa}{b} - 1\right) - h\left(\frac{\rho\sigma xi}{b}\right)\right]w'(h) - \left[\frac{x^2\sigma^2h}{4b^2} + \frac{ix\sigma^2}{4b^2}\right]w(h) = 0.$$
(3.16)

The ODE (3.16) has a general solution of the form [10],

$$w(h) = e^{(\sqrt{1-\rho^2}-\rho i)\frac{\sigma_x}{2b}h} \left[ C_1 \Phi(a^*, b^*, \frac{h}{\zeta}) + C_2 h^{1-b^*} \Phi(a^* - b^* + 1, 2 - b^*, \frac{l}{\zeta}) \right],$$
(3.17)

where

$$a^{*} = \frac{(\sqrt{\rho^{2} - 1} + \rho)\frac{b^{*}}{2} + \frac{\sigma}{4b}}{\sqrt{\rho^{2} - 1}}$$
$$b^{*} = 1 - \frac{\kappa}{b},$$

and

$$\zeta = \frac{-b}{\sigma x \sqrt{1 - \rho^2}}.$$

Here  $C_1$  and  $C_2$  are constants to be determined from the boundary conditions.  $\Phi(a, b, z)$  is the degenerated hypergeometric function which has the following Kummer's series expansion

$$\Phi(a, b, z) = 1 + \sum_{k=1}^{\infty} \frac{(a)_k z^k}{(b)_k k!},$$

where

$$(a)_k = a(a+1)\cdots(a+k-1).$$

If we let  $C_1 = 1$  and  $C_2 = 0$  in (3.17) then a particular solution for (3.16) is

$$w(h) = e^{(\sqrt{1-\rho^2}-\rho i)\frac{\sigma x}{2b}h} \left[\Phi(a^*, b^*, \frac{h}{\zeta})\right].$$

Using the transformation (3.14), Wong and Lo [11] show that a particular solution for (3.13) is

$$U(h) = \frac{2bh}{\sigma^2} \frac{(\sqrt{1-\rho^2}-\rho i)\frac{\sigma x}{2b}\Phi(a^*,b^*,\frac{h}{\zeta}) + \frac{a^*}{b^*\zeta}\Phi(a^*+1,b^*+1,\frac{h}{\zeta})}{\Phi(a^*,b^*,\frac{h}{\zeta})},$$

which can be used to obtain the general solution for (3.13) as follows

$$\hat{D}(h) = U(h) + \frac{\frac{\Phi^2(a^*, b^*, \frac{1}{\zeta})}{\Phi^2(a^*, b^*, \frac{1}{\zeta})} h^{\frac{\kappa}{b}} e^{-2(\sqrt{1-\rho^2})\frac{\sigma x}{2b}(h-1)}}{-\frac{1}{U(1)} + \frac{\sigma^2}{2b} \int_1^h \frac{\Phi^2(a^*, b^*, \frac{1}{\zeta})}{\Phi^2(a^*, b^*, \frac{1}{\zeta})} \eta^{\frac{\kappa}{b} - 1} e^{-2(\sqrt{1-\rho^2})\frac{\sigma x}{2b}(\eta-1)} d\eta}.$$
(3.18)

We now consider the final ordinary differential equation (3.7). Substituting (3.18) and (3.10) in (3.7), we have

$$B_t(t,T) = (\lambda^S m - ba)ixe^{-b(T-t)} - \kappa\theta D(t,T) - \lambda^S \int_{\Re} [e^{ixy} - 1]\phi_Y(y)dy - \lambda^v \int_{\Re} [e^{zD} - 1]\phi_Z(z)dz.$$

Integrating both sides of the above equation and invoking the condition B(T,T) = 0, we obtain

$$B(t,T) = \left(\frac{\lambda^{S}m}{b} - a\right) ix(e^{-b(T-t)} - 1) - \kappa\theta \int_{t}^{T} D(s,T)ds$$
$$+ (T-t)\lambda^{S} \int_{\Re} [e^{ixy} - 1]\phi_{Y}(y)dy$$
$$+ (T-t)\lambda^{v} \int_{\Re} [e^{zD} - 1]\phi_{Z}(z)dz.$$
(3.19)

We can conclude that the characteristic function of the mean reverting process (2.3) with stochastic volatility (2.2) is

$$f(x:t,l,v) = e^{B(t,T) + C(t,T)x + D(t,T)v + ixl},$$

where B(t,T), C(t,T) and D(t,T) are as given in the Lemma.

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## 4 A formula for European Option Pricing

Let C denote the price at time t of a European style call option on the current price of the underlying asset  $S_t$  with strike price K and expiration time T.

The terminal payoff of a European call option on the underlying stock price  $S_t$  with strike price K is

$$\max(S_T - K, 0).$$

This means that the holder will exercise his right only if  $S_T > K$  and then his gain is  $S_T - K$ . Otherwise, if  $S_T \leq K$ , then the holder will buy the underlying asset from the market and the value of the option is zero.

Assuming the risk-free interest rate r is constant over the lifetime of the option, the price of the European call at time t is equal to the discounted conditional expected payoff

$$C(t, S_T) = e^{-r(T-t)} E_{\mathcal{M}}[\max(S_T - K, 0) | \mathcal{F}_t].$$

Assume that t = 0 and we define  $L_T = \ln S_T$  and  $k = \ln K$ . Moreover, we express the call price option  $C(0, S_T)$  as a function of the log of the strike price K rather than the terminal log asset price  $S_T$ . The initial call value  $C_T(k)$  is related to the risk-neutral density  $q_T(l)$  by

$$C_T(k) = e^{-rT} \int_k^\infty (e^l - e^k) q_T(l) dl,$$
(4.1)

where  $q_T(l)$  is the density function of the random variable  $L_T$ . It was mentioned by Carr and Madan [12] that  $C_T(k)$  is not square integrable. To obtain a square integrable function, they introduced the modified call price function  $c_T(k)$  defined by

$$c_T(k) = e^{\alpha k} C_T(k) \tag{4.2}$$

for some constant  $\alpha > 0$  that makes  $c_T(k)$  is square integrable in k over the entire real line and a good choice of  $\alpha$  is that one fourth of the upper bound  $E[S_T^{\alpha+1}] < \infty$ . Consider the Fourier transform of  $c_T(k)$ 

$$\begin{split} \psi_{T}(u) &= \int_{-\infty}^{\infty} e^{iuk} c_{T}(k) dk \\ &= \int_{-\infty}^{\infty} e^{iuk} \int_{k}^{\infty} e^{\alpha k} e^{-rT} (e^{l} - e^{k}) q_{T}(l) dl dk \\ &= \int_{-\infty}^{\infty} e^{-rT} q_{T}(l) \int_{-\infty}^{l} (e^{l+\alpha k} - e^{(1+\alpha)k}) e^{iuk} dk dl \\ &= \int_{-\infty}^{\infty} e^{-rT} q_{T}(l) \left[ \frac{e^{(\alpha+1+iu)l}}{\alpha+iu} - \frac{e^{(\alpha+1+iu)l}}{\alpha+iu+1} \right] dl \\ &= e^{-rT} \int_{-\infty}^{\infty} \left[ \frac{(\alpha+iu)e^{(\alpha+1+iu)l} + e^{(\alpha+1+iu)l} - (\alpha+iu)e^{(\alpha+1+iu)l}}{(\alpha+iu)(\alpha+iu+1)} \right] q_{T}(l) dl \\ &= e^{-rT} \int_{-\infty}^{\infty} \left[ \frac{e^{(\alpha+1+iu)l}}{\alpha^{2} + 2\alpha i u - u^{2} + \alpha + iu} \right] q_{T}(l) dl \end{split}$$

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$$= \frac{e^{-rT}}{\alpha^2 + \alpha - u^2 + i(2\alpha + 1)u} \int_{-\infty}^{\infty} e^{(\alpha + 1 + iu)l} q_T(l) dl$$
  
$$= \frac{e^{-rT}}{\alpha^2 + \alpha - u^2 + i(2\alpha + 1)u} \int_{-\infty}^{\infty} e^{i(u - (\alpha + 1)i)l} q_T(l) dl$$
  
$$= \frac{e^{-rT} f(x = u - (\alpha + 1)i : t, l, v)}{\alpha^2 + \alpha - u^2 + i(2\alpha + 1)u}$$

where f is the characteristic function defined in Lemma 3.1. Hence, the European call prices at time t = 0 with strike price  $k = \ln K$  can then be numerically obtained by using the inverse transform:

$$C_T(k) = \frac{e^{-\alpha k}}{2\pi} \int_{-\infty}^{\infty} e^{-iuk} \psi_T(u) du$$
  
=  $\frac{e^{-\alpha k}}{\pi} \int_0^{\infty} e^{-iuk} \frac{e^{-rT} f(x = u - (\alpha + 1)i : t, l, v)}{\alpha^2 + \alpha - u^2 + i(2\alpha + 1)u} du.$  (4.3)

Integration (4.3) is a direct Fourier transform and lends itself to an application of the Fast Fourier Transform (FFT).

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