# On the Extended Null Scrolls in Minkowski 3-space 

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#### Abstract

In this paper, we defined an extended null scroll in Minkowski 3space $\mathbb{R}_{1}^{3}$ which is obtained by a null line moving with (proper) null frame along a null curve. We proved the well-known theorem due to Bonnet in the 3 -dimensional Euclidean space for an extended null scroll. We calculated the geodesic and normal curvature of a curve on the extended null scroll.


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## 1 Introduction

The Classification of surfaces and curves in Euclidean or Non-Euclidean spaces has been of particular interest for geometers. Many interesting results in NonEuclidean spaces have neen obtained by many mathematicians. This subject have been studied by many researcher $[1,2,3,4,5]$.

In this study, we have done a study about null scrolls in Minkowski 3-sapce. Let $\mathbb{R}_{1}^{3}$ be a Minkowski 3-space with the natural Lorentz metric

$$
\langle\cdot, \cdot\rangle=-d x_{1}^{2}+d x_{2}^{2}+d x_{3}^{2}
$$

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in terms of natural coordinates. The vector product operation of $\mathbb{R}_{1}^{3}$ is defined by

$$
X \Lambda Y=\left(x_{3} y_{2}-x_{2} y_{3}, x_{3} y_{1}-x_{1} y_{3}, x_{1} y_{2}-x_{2} y_{1}\right)
$$

for $X=\left(x_{1}, x_{2}, x_{3}\right), Y=\left(y_{1}, y_{2}, y_{3}\right) \in \mathbb{R}_{1}^{3}$.
The norm of $\vec{v} \in \mathbb{R}_{1}^{3}$ is denoted by $\|\vec{v}\|$ and defined as

$$
\|\vec{v}\|=\sqrt{|\langle\vec{v}, \vec{v}\rangle|} .
$$

Since $\langle\cdot, \cdot\rangle$ is an indefinite metric, recall that a vector $\vec{v} \in \mathbb{R}_{1}^{3}$ can have one of three causal characters: it can be spacelike if $\langle\vec{v}, \vec{v}\rangle>0$ or $\vec{v}=0$, timelike if $\langle\vec{v}, \vec{v}\rangle<0$ and null (lightlike) if $\langle\vec{v}, \vec{v}\rangle=0$ for $\vec{v} \neq 0$. Similarly, an arbitrary curve $\alpha=\alpha(t)$ in $\mathbb{R}_{1}^{3}$ can locally be spacelike, timelike or null (lightlike), if all of its velocity vectors $\alpha^{\prime}(t)$ are respectively spacelike, timelike or null (lightlike) [6].

Definition 1.1. ([7]) A surface in a Minkowski 3-space is called a timelike surface if the induced metric on the surface is a Lorentz metric, i.e., the normal on the surface is a spacelike vector.

Lemma 1.2. ([8]) In the Minkowski 3-space $\mathbb{R}_{1}^{3}$, the following properties are satisfied:
(i) Two timelike vectors are never orthogonal.
(ii) Two null vectors are orthogonal if and only if they are linearly dependent.
(iii) A timelike vector is never orthogonal to a null (lightlike) vector.

A basis $F=\{\vec{X}, \vec{Y}, \vec{Z}\}$ of $\mathbb{R}_{1}^{3}$ is called a (proper) null frame if it satisfies the following conditions:

$$
\begin{aligned}
& \langle\vec{X}, \vec{X}\rangle=\langle\vec{Y}, \vec{Y}\rangle=0, \quad\langle\vec{X}, \vec{Y}\rangle=-1 \\
& \langle\vec{X}, \vec{Z}\rangle=\langle\vec{Y}, \vec{Z}\rangle=0, \quad\langle\vec{Z}, \vec{Z}\rangle=1
\end{aligned}
$$

$\operatorname{det}(X, Y, Z)=1,[9]$.
Let $\alpha=\alpha(t)$ be a null curve in $\mathbb{R}_{1}^{3}$, namely, a smooth curve whose tangent vectors $\alpha^{\prime}(t), \forall t \in I$ are null. For a given smooth positive function $k_{0}=k_{0}(t)$ let us put $X=X(t)=k_{0}^{-1} \alpha^{\prime}$. Then $X$ is a null vector field along $\alpha$. Moreover, there exists a null vector field $Y=Y(t)$ along $\alpha$ satisfying $\langle X, Y\rangle=-1$. Here if we put $Z=X \Lambda Y$ then we can obtain a (proper) null frame field $F=\{X, Y, Z\}$ along $\alpha$. In this case the pair $(\alpha, F)$ is said to be a (proper) framed null curve. A framed null curve ( $\alpha, F$ ) satisfies the following, so called the Frenet equations:

$$
\begin{align*}
X^{\prime}(t) & =k_{1}(t) X(t)+k_{2}(t) Z(t), \\
Y^{\prime}(t) & =-k_{1}(t) Y(t)+k_{3}(t) Z(t),  \tag{1.1}\\
Z^{\prime}(t) & =k_{3}(t) X(t)+k_{2}(t) Y(t),
\end{align*}
$$

where $k_{i}=k_{i}(t), i=1,2,3$ are smooth functions defined by

$$
k_{1}=-\left\langle X^{\prime}, Y\right\rangle, \quad k_{2}=\left\langle X^{\prime}, Z\right\rangle, \quad k_{3}=\left\langle Y^{\prime}, Z\right\rangle
$$

The function $k_{i}$ is called an $i-t h$ curvature of the framed null curve. It follows from the fundamental theorem of ordinary differential equations that a framed null curve $(\alpha, F)=(\alpha(t), F(t))$ is uniquely determined by the functions $k_{0}(>0), k_{1}, k_{2}, k_{3}$ and the initial condition. The functions $k_{2}$ and $k_{3}$ are called the curvature and torsion of $\alpha$, respectively.

A framed null curve $(\alpha, F)$ with $k_{0}=1$ and $k_{1}=0$ is called a Cartan framed null curve and the frame field $F$ is called a Cartan frame.

We call the vector fields $X, Y, Z$ a tangent vector field, a binormal vector field and a (principal) normal vector field of $\alpha$, respectively.

Note that $\alpha$ is called null geodesic if $k_{2}=0[9,10]$.
Lemma 1.3. ([11]) Assume that $\alpha(t)$ is a null curve in $\mathbb{R}_{1}^{3}$ and $\{X, Y, Z\}$ be its (proper) null frame field. Then

$$
X \Lambda Y=Z, \quad Y \Lambda Z=-Y, \quad X \Lambda Z=X
$$

## 2 Null Scrolls in $\mathbb{R}_{1}^{3}$

Let $(\alpha, F)=(\alpha(t), F(t))$ be a null curve with frame $F=\{X, Y, Z\}$. A ruled surface is a surface swept out by a straight line Y moving along a curve $\alpha$. The various positions of the generating line Y are called the rulings of the surface. Such a surface, thus has a parametrization in ruled form as follows:

$$
\begin{gathered}
\Psi: U \longrightarrow \mathbb{R}_{1}^{3} \\
\Psi(t, v)=\alpha(t)+v Y(t) ; \quad \forall(t, v) \in U
\end{gathered}
$$

is called a null scroll and denoted by $M$. We call $\alpha$ to be the base curve and $Y$ to be the director curve. If the tangent plane is constant along each ruling, then the ruled surface is called a developable surface. The remaining ruled surfaces are called skew surfaces. One can see that $M$ is a timelike surface. Furthermore, for a Cartan framed null curve $\alpha$ with Cartan frame $F=\{X, Y, Z\}$ the ruled surface is called a B-scroll [11].

Now consider a ruled surface in $\mathbb{R}_{1}^{3}$ generated by a null generator $\vec{L}(t)$ moving with (proper) null frame of a null curve $\alpha=\vec{\alpha}(t)$, i.e.,

$$
\begin{equation*}
\vec{L}(t)=\ell_{1}(t) X(t)+\ell_{2}(t) Y(t)+\ell_{3}(t) Z(t) \tag{2.1}
\end{equation*}
$$

where the components $\ell_{i}=\ell_{i}(t), \ell_{2} \neq 0(i=1,2,3)$ are scalar functions of the parameter of the null curve $\alpha=\vec{\alpha}(t)$. Thus if $\vec{L}$ moves with (proper) null frame, the constructed ruled surface is given by the following parametrization

$$
\Psi(t, v)=\vec{\alpha}(t)+v \vec{L}(t), \quad(t, v) \in U \subset \mathbb{R}^{2}
$$

$$
\begin{equation*}
<\vec{L}(t), \vec{L}(t)>=\ell_{3}^{2}(t)-2 \ell_{1}(t) \ell_{2}(t)=0 \tag{2.2}
\end{equation*}
$$

This ruled surface is called an extended null scroll.
From (2.1) and using (1.1), we obtain

$$
\begin{equation*}
\overrightarrow{L^{\prime}}(t)=\left(\ell_{1}^{\prime}+\ell_{1} k_{1}+\ell_{3} k_{3}\right) \vec{X}+\left(\ell_{2}^{\prime}-\ell_{2} k_{1}+\ell_{3} k_{2}\right) \vec{Y}+\left(\ell_{3}^{\prime}+\ell_{1} k_{2}+\ell_{2} k_{3}\right) \vec{Z} \tag{2.3}
\end{equation*}
$$

It obvious that the vector $\overrightarrow{L^{\prime}}(t)$ is a spacelike vector or null vector and in the second case it is linearly dependent with the generator.

The assumption $\overrightarrow{L^{\prime}}(t) \neq 0$, is usually expressed by saying that the ruled surface $M$ is a noncylindrical.

From (2.2), one can obtain the first fundamental quantities of the extension

$$
\begin{equation*}
g_{11}=-2 k_{0} v\left(\ell_{2}^{\prime}-\ell_{2} k_{1}+\ell_{3} k_{2}\right)+v^{2}\left\|L^{\prime}\right\|^{2}, \quad g_{12}=-k_{0} \ell_{2}, \quad g_{22}=0 \tag{2.4}
\end{equation*}
$$

Thus the induced metric on the extended null scroll is a Lorentz metric. Therefore the extended null scroll is a timelike ruled surface.

The unit normal vector field $\vec{n}=\vec{n}(t, v)$ on the extended null scroll in $\mathbb{R}_{1}^{3}$ is

$$
\begin{equation*}
\vec{n}=\frac{\alpha^{\prime}(t) \wedge \vec{L}(t)+v \overrightarrow{L^{\prime}}(t) \wedge \vec{L}(t)}{\left\|\alpha^{\prime}(t) \wedge \vec{L}(t)+v \overrightarrow{L^{\prime}}(t) \wedge \vec{L}(t)\right\|} \tag{2.5}
\end{equation*}
$$

Thus, from (2.5) the unit normal vector to the surface $M$ at the point $(\mathrm{t}, 0)$ is

$$
\begin{equation*}
\vec{n}(t, 0)=\frac{\ell_{3} X+\ell_{2} Z}{\left|\ell_{2}\right|}, \quad \ell_{2} \neq 0 \tag{2.6}
\end{equation*}
$$

Definition 2.1. ([1]) If there exists a common perpendicular to two preceding rulings in the skew surface, the foot of the common perpendicular on the main ruling is called a central point. The locus of the central points is called the curve of striction.

Using (1.1) and (2.3), it is easy to see that the parametrization of the striction curve on an extended null scroll (2.2) is given by

$$
\begin{equation*}
\vec{\beta}(t)=\vec{\alpha}(t)+\frac{k_{0}\left(\ell_{2}^{\prime}-\ell_{2} k_{1}+\ell_{3} k_{2}\right)}{\left\|\overrightarrow{L^{\prime}}\right\|^{2}} \vec{L}(t) \tag{2.7}
\end{equation*}
$$

From (2.6), it follows that the base curve of the extended null scroll (2.2) to be a geodesic curve ( $\vec{n}(t, 0)=\vec{Z}$ ) if $\ell_{2} \neq 0$ and $\ell_{3}=0$.

Theorem 2.2. Let $M$ be an extended null scroll. Then the curve of striction is a timelike curve in an extended null scroll M.

Proof. If we use the equation (2.6), we can show easily that the tangent vector field of the curve of striction is a timelike vector field.

Thus we have the Bonnet's theorem for the extended null scroll which is formulated as follows:

Theorem 2.3. If a timelike curve on an extended null scroll in $\mathbb{R}_{1}^{3}$ has two of the following properties, it has the third also
(i) it is a geodesic $\left(k_{2}=0\right)$ and $k_{1}=0$,
(ii) it cut the rulings at a constant angle ( $\ell_{2}=$ const.),
(iii) it is a striction curve.

Now we examine the extended null scroll for which the striction curve is the base curve, i.e.,

$$
\begin{gather*}
M^{s}: \Psi^{s}(t, v)=\vec{\alpha}(t)+v \vec{L}(t) \\
\ell_{3}^{2}-2 \ell_{1} \ell_{2}=0, \quad \ell_{2}^{\prime}-\ell_{2} k_{1}+\ell_{3} k_{2}=0 . \tag{2.8}
\end{gather*}
$$

Corollary 2.4. Let the curve of striction be the base curve of an extended null scroll $M$. Then the vector $\overrightarrow{L^{\prime}}(t)$ is a spacelike vector.

Proof. From (2.3), using (2.8), it can be seen easily.
Let us consider an extended null scroll $M^{s}$. Then since $\left\langle\vec{L}(t), \overrightarrow{L^{\prime}}(t)\right\rangle=0$ and $\left\langle\overrightarrow{\alpha^{\prime}}, \overrightarrow{L^{\prime}}(t)\right\rangle=0$, we can write

$$
\overrightarrow{\alpha^{\prime}}(t) \wedge \vec{L}(t)=\lambda \overrightarrow{L^{\prime}}(t)
$$

where

$$
\begin{equation*}
\lambda=\lambda(t)=-\frac{\operatorname{det}\left(\overrightarrow{\alpha^{\prime}}(t), \vec{L}(t), \overrightarrow{L^{\prime}}(t)\right)}{\left\|\overrightarrow{L^{\prime}}(t)\right\|^{2}} . \tag{2.9}
\end{equation*}
$$

The function $\lambda=\lambda(t)$ is called the distribution parameter of the extended null scroll $M^{s}$. In more explicitly using (1.1), (2.1) and (2.3), we get

$$
\begin{equation*}
\lambda(t)=-\frac{k_{0} \ell_{2}\left(\ell_{1} k_{2}+\ell_{2} k_{3}+\ell_{3}^{\prime}\right)}{\left\|\overrightarrow{L^{\prime}}(t)\right\|^{2}} . \tag{2.10}
\end{equation*}
$$

The normal vector field on $M^{s}$ takes the form

$$
\begin{equation*}
\vec{N}=\lambda \overrightarrow{L^{\prime}}(t)+v \overrightarrow{L^{\prime}}(t) \wedge \vec{L}(t) \tag{2.11}
\end{equation*}
$$

Thus, from (2.10) we have

$$
\begin{equation*}
\ell_{2}^{2}=\frac{\lambda^{2}}{k_{0}^{2}}\left\|\overrightarrow{L^{\prime}}(t)\right\|^{2} \tag{2.12}
\end{equation*}
$$

Hence, we have

Corollary 2.5. The singular points on the extended null scroll $M^{s}$ are the points for which $\lambda=0$. Since $\left\|\overrightarrow{L^{\prime}}\right\| \neq 0$, i.e., $\ell_{1} k_{2}+\ell_{2} k_{3}+\ell_{3}^{\prime} \neq 0$, the singular points are given $\ell_{2}=0$.

Let $M^{s}$ be an extended null scroll for which the striction curve is the base curve in $\mathbb{R}_{1}^{3}$. The unit normal vector to the extended null scroll $M^{s}$ at $(t, v)$ is given from (2.11) and (2.12) by

$$
\vec{n}(t, v)=k_{0} \vec{\ell}+k_{0} \frac{v}{\lambda} \vec{\ell} \wedge \vec{L},
$$

where $\vec{\ell}=\frac{\overrightarrow{L^{\prime}}(t)}{\left\|\overrightarrow{L^{\prime}}(t)\right\|}$.
For a regular patch on $M^{s}(\lambda \neq 0)$, it is easy to see that the normal along the striction curve on $M^{s}$ is given by

$$
\vec{n}_{0}(t, 0)=k_{0} \vec{\ell} .
$$

Since $\vec{n}_{0}$ is a unit spacelike vector and $\vec{n}$ is unit spacelike vector, thus if $\theta$ is the angle of rotation from the normal $\vec{n}_{0}$ to the normal $\vec{n}$ we get

$$
\sin \theta=\left\|\vec{n}_{0} \wedge \vec{n}\right\|=\left\|k_{0} \vec{\ell}(t)+v \frac{k_{0}}{\lambda} \vec{\ell}(t) \wedge \vec{L}(t) \wedge \vec{\ell}(t)\right\| .
$$

Routine calculation, one can obtain $\theta=0$. Thus we have without loose of generality following theorem which is similar to Chasles theorem for the extended null scroll in $\mathbb{R}_{1}^{3}$.

Theorem 2.6. For the extended null scroll $M^{s}$ in $\mathbb{R}_{1}^{3}$, the normal vector $\vec{n}$ at a point of a ruling and the normal vector $\vec{n}_{0}$ at the striction point of this ruling are parallel.

## 3 Extended Null Scrolls with Constant Parameter of Distribution

From (2.10), one can see that an extended null scroll $M^{s}$ with a constant distribution parameter satisfies the following differential equation:

$$
-k_{0} \ell_{2}\left(\ell_{1} \ell_{2}+\ell_{2} k_{3}+\ell_{3}^{\prime}=c\left\langle\overrightarrow{L^{\prime}(t)}, \overrightarrow{L^{\prime}(t)}\right\rangle,\right.
$$

where $c$ is constant. If we consider (2.8), we get easily,

$$
\begin{align*}
\ell_{2} & =\int\left(k_{1} \ell_{2}-k_{2} \ell_{3}\right) d t+c_{1} \\
\ell_{3} & =-\int\left(\frac{{ }^{c}\left\langle\overrightarrow{L^{\prime}(t)}, \overrightarrow{L^{\prime}(t)}\right\rangle}{k_{0} \ell_{2}}+\frac{\ell_{3}^{2}}{2}+k_{3} \ell_{2}\right) d t+c_{2}  \tag{3.1}\\
\ell_{3}^{2}-2 \ell_{1} \ell_{2} & =0 .
\end{align*}
$$

The first equation of (3.1) is an integral equation for the unknown $\ell_{1}=\ell_{1}(t)$. Therefore if $\ell_{2}=\ell_{2}(t)$ and $\ell_{3}=\ell_{3}(t)$ are given we get $\ell_{1}=\ell_{1}(t)$.

Theorem 3.1. The range of existence of a one parametric extended null scrolls $\left\{M^{s}\right\}$ with constant parameter of distribution comprises within two arbitrary functions of one variable.

Developable null scrolls are special class of the ruled surfaces which is described $k_{1}=k_{3}=0[11]$.

Definition 3.2. ([11]) Let $M^{s}$ be an extended null scroll in $\mathbb{R}_{1}^{3}$. If there exists a curve which makes constant angle with one of the rulings, this curve is called a pseudo-orthogonal trajectory of $M^{s}$.

Theorem 3.3. Let $M$ be an extended null scroll in $\mathbb{R}_{1}^{3}$. The shortest distance between two ruling is measured only on the curve of striction which is one of the pseudo-orthogonal trajectories.

Now we calculate the geodesic curvature and normal curvature of an extended null scroll $M^{s}$ in $\mathbb{R}_{1}^{3}$.

Let $\vec{\gamma}=\vec{\gamma}(t)$ be a curve on the extended null scroll $M^{s}$. Then it can be represented in the form

$$
\vec{\gamma}(t)=\vec{\alpha}(t)+v(t) \vec{L}(t)
$$

Using (1.1), it is easy to see that the unit tangent vector along the curve $\vec{\gamma}=\vec{\gamma}(t)$ is

$$
\vec{\gamma}_{0}^{\prime}(t)=\frac{\gamma^{\prime}(t)}{\left\|\gamma^{\prime}(t)\right\|}=\frac{\xi_{1} X+\xi_{2} Y+\xi_{3} Z}{\sqrt{\left|\xi_{3}^{2}-2 \xi_{1} \xi_{2}\right|}}
$$

where,

$$
\begin{aligned}
\xi_{1} & =k_{0} v^{\prime} \ell_{1}+v \eta_{1} \\
\xi_{2} & =\eta_{2} \\
\xi_{3} & =v^{\prime} \ell_{3}+v \eta_{3} \\
\eta_{1} & =\ell_{1}^{\prime}+\ell_{1} k_{1}+\ell_{3} k_{3} \\
\eta_{2} & =v^{\prime} \ell_{2} \\
\eta_{3} & =\ell_{3}^{\prime}+\ell_{1} k_{2}+\ell_{2} k_{3}
\end{aligned}
$$

Therefore, one can see that

$$
\begin{aligned}
\vec{\gamma}_{0}^{\prime \prime}= & \frac{1}{R^{\frac{1}{2}}}\left[\left(\xi_{1}^{\prime}+\xi_{1} k_{1}+\xi_{3} k_{3}\right) X+\left(\xi_{2}^{\prime}-\xi_{2} k_{1}+\xi_{3} k_{2}\right) Y+\left(\xi_{3}^{\prime}+\xi_{1} k_{2}+\xi_{2} k_{3}\right) Z\right] \\
& -\frac{R^{\prime}}{2 R^{\frac{3}{2}}}\left(\xi_{1} X+\xi_{2} Y+\xi_{3} Z\right)
\end{aligned}
$$

where $R=\left|\xi_{3}^{2}-2 \xi_{1} \xi_{2}\right|$.

Using (2.5), the unit normal vector field on the extended null scroll $M^{s}$ along the curve $(v=v(t))$ is

$$
\vec{n}(t, v(t))=\frac{\left[k_{0} \ell_{3}+v\left(\ell_{3} \eta_{1}-\ell_{1} \eta_{3}\right)\right] X+\left(v \ell_{2} \eta_{3}\right) Y+\left(k_{0} \ell_{2}+v \ell_{2} \eta_{1}\right) Z}{\sqrt{\left(k_{0} \ell_{2}+v \ell_{2} \eta_{1}\right)^{2}-2 v \ell_{2} \eta_{3}\left[k_{0} \ell_{3}+v\left(\ell_{3} \eta_{1}-\ell_{1} \eta_{3}\right)\right]} .} .
$$

The geodesic curvature of the curve $\vec{\gamma}=\vec{\gamma}(t)$ is given by

$$
k_{g}=\frac{1}{2\|\vec{n}(t, v(t))\| R^{\frac{3}{2}}}\left|\begin{array}{ccc}
\xi_{1} & \xi_{2} & \xi_{3} \\
2 R\left(\xi_{1}^{\prime}+\xi_{1} k_{1}\right. & 2 R\left(\xi_{2}^{\prime}-\xi_{2} k_{1}\right. & 2 R\left(\xi_{3}^{\prime}+\xi_{1} k_{2}\right. \\
\left.+\xi_{3} k_{3}\right)-R^{\prime} \xi_{1} & \left.+\xi_{3} k_{2}\right)-R^{\prime} \xi_{2} & \left.+\xi_{2} k_{3}\right)-R^{\prime} \xi_{3} \\
k_{0} \ell_{3} \\
+v\left(\ell_{3} \eta_{1}-\ell_{1} \eta_{3}\right) & v \ell_{2} \eta_{3} & k_{0} \ell_{2}+v \ell_{2} \eta_{1}
\end{array}\right| .
$$

Then, the normal curvature of a curve $\overrightarrow{\gamma_{0}}=\overrightarrow{\gamma_{0}}(t)$ is given by

$$
\begin{aligned}
k_{n}= & \frac{1}{2\|\vec{n}(t, v(t))\| R^{\frac{3}{2}}}\left\{-\left[2 R\left(\xi_{1}^{\prime}+\xi_{1} k_{1}+\xi_{3} k_{3}\right)-R^{\prime} \xi_{1}\right]\left(v \ell_{2} \eta_{3}\right)\right. \\
& -\left[2 R\left(\xi_{2}^{\prime}-\xi_{2} k_{1}+\xi_{3} k_{2}\right)-R^{\prime} \xi_{2}\right]\left(k_{0} \ell_{3}+v\left(\ell_{3} \eta_{1}-\ell_{1} \eta_{3}\right)\right) \\
& \left.+\left[2 R\left(\xi_{3}^{\prime}+\xi_{1} k_{2}+\xi_{2} k_{3}\right)-R^{\prime} \xi_{3}\right]\left(k_{0} \ell_{2}+v \ell_{2} \eta_{1}\right)\right\} .
\end{aligned}
$$

Corollary 3.4. From (3.2) and (3.3) the geodesic curvature $k_{g}$ and the normal curvature $k_{n}$ of a base curve on the extended null scroll $M^{s}$ in $\mathbb{R}_{1}^{3}$ are not define.

The results in the study are confirmed by the following example:
Example 3.5. $\varphi(t, v)=(t+v, \cos t+v \cos t, \sin t+v \sin t)$ is a null scroll where $\alpha(t)=(t, \cos t, \sin t)$ is a null base curve and $\vec{L}(t)=(1, \cos t, \sin t)$ is a null generator. The striction curve is $\bar{\alpha}(t)=\alpha(t)-\vec{L}(t)$. The distribution parameter is $\lambda=-1$ and it is nondevelopable null scroll.


Fig1.

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